



## Enhanced Optical Response of Alumino-borophosphate Glasses Embedded with Gold Nanoparticles for Photonic Application

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### ABSTRACT

The demand for efficient, energy-saving, and environmentally friendly lighting devices remains a critical pursuit in materials science. In response to this need, we have developed a novel series of gold nanoparticles embedded in alumino-borophosphate glasses, using the conventional melt quenching technique. These glasses were subjected to both experimental and theoretical analyses to evaluate their potential for lasing and light-emitting applications. X-ray diffraction (XRD) confirmed the amorphous structure of the glass, while high-resolution transmission electron microscopy (HRTEM) images revealed spherical gold nanoparticles, averaging 15.01 nm in size, embedded within the glass matrix. The experimental absorption and photoluminescence spectra were further validated through Judd–Ofelt intensity and radiative parameter calculations. Notably, the incorporation of gold nanoparticles significantly enhanced the optical properties of the glass. In particular, the glass sample containing 0.7 mol% of gold nanoparticles exhibited a high branching ratio (72%), stimulated emission cross-section ( $30.24 \times 10^{-20} \text{ cm}^2$ ), quantum efficiency (95%), and exceptional color purity (89%). The strong optical response achieved in this study underscores the suitability of these fabricated glasses for the development of an innovative solid-state red laser operating at 611 nm.

**Keywords:** Alumino-borophosphate glasses; Gold nanoparticles; Optical properties; Judd Ofelt parameters.

### INTRODUCTION

The promising features of this glass host, which allows for tailoring the spectroscopic properties of rare earth-doped boro-phosphate glasses through surface plasmon resonance (SPR) by embedding metallic nanoparticles (NPs), present a challenging task. Achieving an optimal concentration of NPs is crucial for enhancing the optical response, making this

approach valuable for photonic applications [1]. Equally, luminescence quenching of rare earth ions is detrimental for fabricating efficient lasers [2-3]. This drawback needs to be overcome, and significant enhancement of spectral features is essential [4-6]. Embedding metallic NPs/ or co-doping by low concentration of rare earth ions have proven to be alternative way to prevail such shortcomings [7-9]. Keeping in view of the

significant and potential advantages of metallic nanoparticles embedded rare earth co-doped glass materials, the proposed research aims to investigate the influence of gold (Au) NPs on the optical response of  $\text{Nd}^{3+}/\text{Er}^{3+}$  Co-doped Alumino boro-phosphate glasses [10-11]. Based on the absorption and emission spectra, the spectroscopic properties such as spontaneous radiative emission probability ( $A_T$ ), radiative lifetime ( $\tau_R$ ), luminescence branching ratio ( $\beta_R$ ), effective bandwidth ( $\Delta\lambda_{\text{eff}}$ ), stimulated emission cross section ( $\sigma_p$ ), gain bandwidth ( $\sigma_p \times \Delta\lambda_{\text{eff}}$ ), and optical gain ( $\sigma_p \times \tau_R$ ) of Au NPs embedded  $\text{Nd}^{3+}/\text{Er}^{3+}$  Co-doped alumino-borophosphate glasses will be evaluated using Judd–Ofelt theory to ascertain their utility for solid state laser and light emitting devices [6-12]. Although numerous studies have explored phosphate glass systems, no definitive conclusions have been reached regarding the optimal composition or the ideal combination of nanoparticles and rare earth ions in these glasses. In this context, our findings may offer valuable insights toward the development of high-performance phosphate glass-based solid-state laser media, which are essential for a range of functional applications [13-14].

## MATERIALS AND METHODS

### Glass Sample Synthesis and Characterization

In this study, glasses were successfully synthesized using the conventional melt quenching technique. A 15-gram mixture was combined in a platinum crucible and subsequently melted at  $900^\circ\text{C}$  for 40 minutes. All raw materials were sourced from Sigma Aldrich Company. Once a suitable viscous mixture was achieved, the melt was annealed

at  $300^\circ\text{C}$  for 3 hours by quenching it onto a stainless-steel mold. The furnace was then gradually cooled down to room temperature. The glass samples produced were visually transparent. They were characterized both thermally and optically using conventional spectroscopic techniques. Absorption analysis was carried out at room temperature using a monochromatic (Monk-Gillieson) spectrophotometer fitted with a photodiode detector, covering a wavelength range of 220 to 2600 nm. To ascertain the phase nature of the glass samples, X-ray diffraction (XRD) patterns were acquired employing a Siemens D5000 X-ray diffractometer.

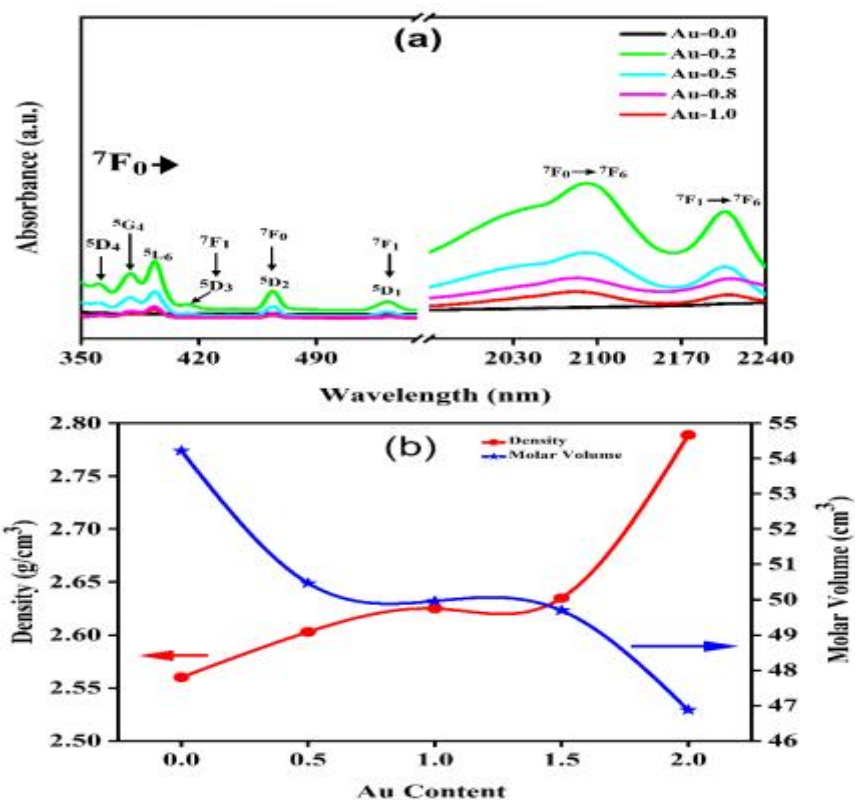
## RESULTS AND DISCUSSION

### 3.1 Physical and Optical properties of the glasses

Table 1 presents the physical and optical properties of the glasses, including density, molar volume, ionic concentration, polar radius, and molar refractivity. Figure 1(a) depicts the absorbance versus wavelength for the glasses in the UV-visible and near-infrared (NIR) regions, recorded over a wavelength range of 350 to 2240 nm, while Figure 1(b) shows the density and molar volume of the prepared glass system. As seen in Figure 1(a), there are approximately seven prominent absorption peaks across the ultraviolet-visible to near-infrared spectrum. The identified transitions in the ultraviolet (UV) region include  ${}^7\text{F}_0 \rightarrow {}^5\text{D}_4$  (360 nm),  ${}^7\text{F}_0 \rightarrow {}^5\text{G}_4$  (379 nm),  ${}^7\text{F}_0 \rightarrow {}^5\text{L}_6$  (394 nm),  ${}^7\text{F}_1 \rightarrow {}^5\text{D}_3$  (414 nm),  ${}^7\text{F}_0 \rightarrow {}^5\text{D}_2$  (464 nm),  ${}^7\text{F}_1 \rightarrow {}^5\text{D}_1$  (532 nm). Additionally, the NIR region shows transitions at  ${}^7\text{F}_0 \rightarrow {}^5\text{F}_6$  (2090 nm), and  ${}^7\text{F}_1 \rightarrow {}^5\text{F}_6$  (2205 nm). These transitions align with those reported in the literature [9,10].

**Table 1:** Physical and optical properties Au glasses

| Parameters                                  | Au-0.0 | Au-0.2 | Au-0.5 | Au-0.8 | Au-1.0 |
|---|--------|--------|--------|--------|--------|
| $\rho$ (g/cm <sup>3</sup> )                 | 2.560  | 2.603  | 2.625  | 2.635  | 2.789  |
| $V_m$ (cm <sup>3</sup> /mol)                | 54.22  | 50.47  | 49.96  | 49.69  | 46.87  |
| $N \times 10^{22}$ (ions cm <sup>-3</sup> ) | 0.664  | 0.556  | 1.220  | 1.831  | 2.573  |
| $R_p$ (Å)                                   | 4.320  | 2.312  | 1.781  | 1.564  | 1.396  |
| $F$ ( $\times 10^{16}$ cm <sup>-2</sup> )   | 7.258  | 11.84  | 19.89  | 26.01  | 32.63  |
| $R_i$ ( $\times 10^{-23}$ Å)                | 7.010  | 6.670  | 5.151  | 4.501  | 4.020  |
| $R$ %                                       | 0.991  | 1.371  | 1.411  | 1.503  | 1.542  |
| $R_M$                                       | 46.11  | 30.02  | 27.57  | 27.23  | 26.91  |
| $\alpha_e$ ( $\times 10^{-23}$ )            | 16.33  | 11.93  | 11.14  | 10.92  | 10.61  |


**Figure 1:** (a) Absorbance as function of wavelength and (b) density and molar volume of the prepared glass system

### Structural properties

#### *X-rays diffraction and Morphological structure*

The XRD analysis shows no sharp peaks, indicating that the samples are amorphous and possess short-range order, as evidenced by a broad hump between 20° and 25° (2 $\theta$ ) in Figure 2. The absence of distinct crystalline peaks, combined with the broad hump,

confirms the amorphous structure of the samples [10]. In addition, the low-resolution TEM image in Figure 3(a) clearly reveals gold nanoparticles (AuNPs) as dark spots. Figure 3(b) provides an HRTEM image along with a Fast Fourier Transform (FFT) diagram, highlighting the lattice fringes of the crystalline structure of the AuNPs. Notably, only the gold nanoparticles are visible,

consistent with earlier studies [15-17]. The lattice fringes of the AuNPs are typically well-defined [18-19]. Furthermore, the TEM image of the selected Au-0.3 glass (Figure 3c)

displays a uniform distribution of nearly spherical nanoparticles within the glass matrix, with an average diameter of approximately 27.65 nm, as shown in Figure 3(d).

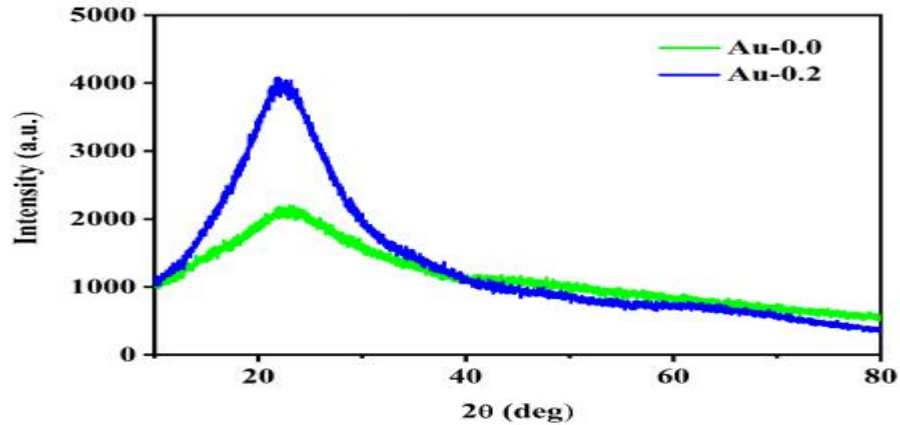


Figure 2: X-rays diffraction of the prepared glass system.

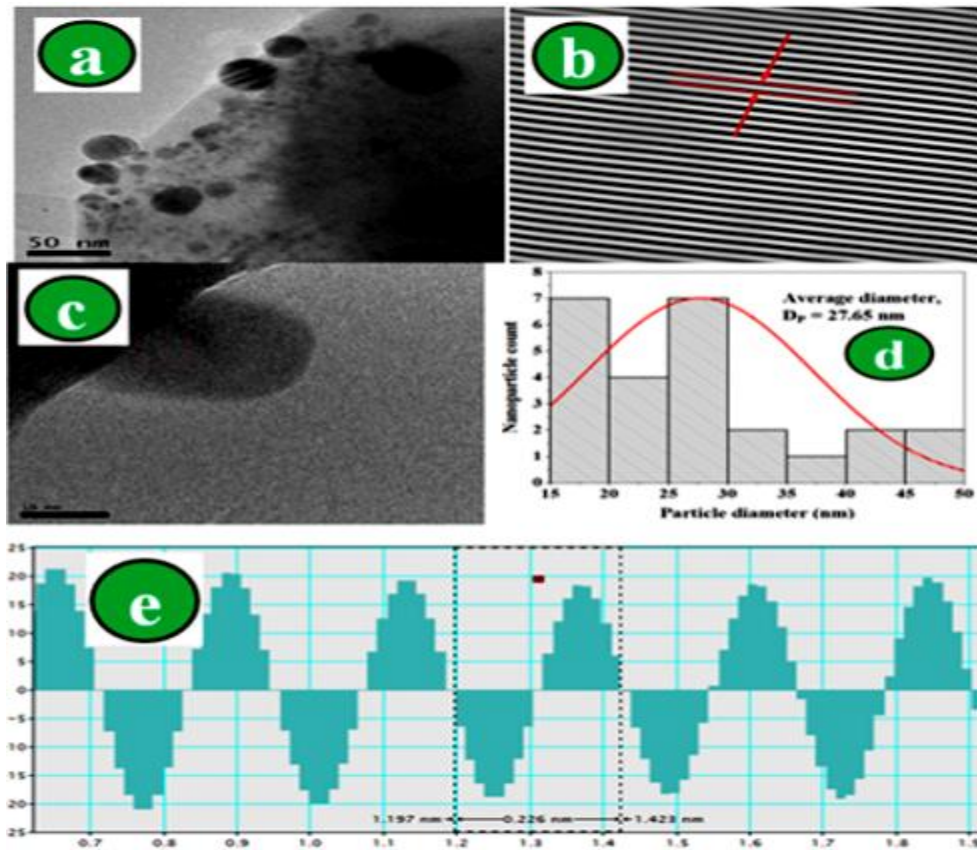


Figure 3: (a) HRTEM micrograph images of spherical AuNPs at different grid location (b) a single AuNPs with lattice fringes, (c) HRTEM image of a single AuNPs (d) average diameter of AuNPs of Au0.50 glass system

## Theoretical Analysis

### Judd-Ofelt analysis

The Judd-Ofelt intensity parameter is an effective tool for analyzing the spectroscopic properties of materials, utilizing either absorption or emission spectra. For certain materials, however, emission spectra are necessary due to the lack of specific transitions [20,21]. The simulated emission

$$\Omega_\lambda = \frac{A_{ed}}{A_{md}} \left\{ \frac{S_{md} \nu_1^3}{e^2 \nu_\lambda^3} \right\} \left\{ \frac{9n^2}{(n^2 + 2)^2} \right\} \left[ \left| \langle \psi J \| U^\lambda \| \psi' J' \rangle \right|^2 \right]^{-1}, \quad (1)$$

where  $n$  is refractive index,  $S_{md}$  is magnetic dipole,  $A_{ed}$  and  $A_{md}$  are electric and magnetic dipole transition probabilities and  $\nu_\lambda$  are transitions wavenumber.

The calculated data of the Judd-Ofelt intensity parameters and others related works are summarized in Table 3. The Au-0.5 glass demonstrates a high value of  $\Omega_2$   $7.55 \times 10^{-20} \text{cm}^2$ .

cross-section is crucial for understanding lasing transitions. A higher  $\Omega_2$  value indicates greater asymmetry in rare earth ions and ligand field effects, while  $\Omega_4$  is influenced by transitional ions, the material's bulk properties, and its rigidity. The absence of the  ${}^5D_0 \rightarrow {}^7F_6$  transition results in  $\Omega_6$  to be zero [9, 21]. The  $\Omega_\lambda$  value is derived from the intensity ratios of transitions and is given by

Conversely, the Au-0.8 glass displays the lowest value of  $\Omega_2$  at  $6.21 \times 10^{-20} \text{cm}^2$  as concentration of AuNPs increases from 0.5 to 1.0g.

**Table 2:** Calculated data of Judd-Ofelt parameters of the studied glasses and other related works.

| Glass code | $\Omega_2$ | $\Omega_4$ | $\Omega_6$ | Trend                            | Reference  |
|------------|------------|------------|------------|----------------------------------|------------|
| Au-0.0     | 6.41       | 2.70       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | This study |
| Au-0.2     | 6.88       | 2.36       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | This study |
| Au-0.5     | 7.55       | 2.42       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | This study |
| Au-0.8     | 6.21       | 2.05       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | This study |
| Au-1.0     | 7.45       | 2.06       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | This study |
| TBZEu0.03  | 3.37       | 1.82       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | [9]        |
| PKMNC      | 5.62       | 2.05       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | [22]       |
| CSEu0.1    | 3.55       | 0          | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | [23]       |
| BGG        | 5.45       | 4.45       | 0          | $\Omega_2 > \Omega_4 > \Omega_6$ | [24]       |

\*<sup>9</sup>Borotellurite, \*<sup>22</sup>Phosphate, \*<sup>23</sup>Calcium silicate, \*<sup>24</sup>Borogermanate.

**Table 3.** Radiative parameter, branching ratio and stimulated emission cross-section of  ${}^5D_0 \rightarrow {}^7F_2$  transition of studied glasses and others related studies.

| Glass code | $A_{rad}, \text{s}^{-1}$ | $\beta_r, \%$ | $\sigma_{em}, (\times 10^{-22} \text{cm}^2)$ | Reference  |
|------------|--------------------------|---------------|--|------------|
| Au-0.2     | 375.00                   | 61.6          | 9.55   | This study |
| BPEu0.25   | 125.49                   | 61.0          | 10.05  | [16]       |
| PAZE20     | 219.00                   | 71.0          | 8.36   | [17]       |
| TBZEu0.03  | 565.10                   | 61.1          | 12.57  | [9]        |
| TZNSAu-4   | 357.00                   | 37            | 0.143  | [25]       |
| TZNSAu-3   | 294.00                   | 45            | 0.273  | [25]       |
| BTBaEu0.05 | 150.00                   | 69            | 17.5   | [26]       |

\*<sup>16</sup>Borophosphate, \*<sup>17</sup>Sodium phosphate, \*<sup>9</sup>Borotellurite, \*<sup>25</sup>Sodium zinc-tellurite, \*<sup>26</sup>Barium-telluraborate



### Radiative parameters

The critical parameters, such as the stimulated emission cross-section, are essential in evaluating the optical excellence of the glass for potential laser applications. A branching ratio exceeding 50% is necessary for lasing,

$$A_{ed} = \left[ \frac{64\pi^4 \nu^3 n(n^2 + 2)^2}{27\hbar(2J+1)} \right] S_{ed} \quad (2)$$

where in equation (2). The branching ratio ( $\beta_r$ ) is the radiative transition probability ratio to sum of radiative transition in all levels. Meanwhile,  $\tau_{cal}$  is defined as a calculated lifetime associated with the sum of all radiative transition. The branching ratio and  $\tau_{cal}$  are given as expressed as follows:

$$\beta_r = \frac{A_{rad}}{\sum A_{rad}} \quad (3)$$

$$\tau_{cal} = \frac{1}{\sum A_{rad}} \quad (5)$$

Furthermore, emission cross-section is given by

$$A_e = \frac{\lambda P^4}{8\pi c n^2 \Delta_{eff}} A(\psi j) \quad (6)$$

where  $\lambda_p$  emission peak,  $n$  is the refractive index, and  $c$  is the speed of light in vacuum.

Table 4 presents the transition probabilities ( $A_{rad}$ ), branching ratios ( $\beta_r$ ), emission cross-section of the excited state, optical bandwidth, and optical gain width for the  ${}^5D_0 \rightarrow {}^7F_2$  transitions of the co-embedded nanoparticles [15,28]. The importance of the branching ratio in laser development highlights the need for higher stimulated emission values in glass transitions [9]. The trend in intensity parameters suggests structural changes in the surrounding rare earth ions, indicating a dominance of short-range influences and covalent characteristics, while  $\Omega_4$  is associated with bulk properties such as rigidity and viscosity [9,15]. A higher intensity of  $\Omega_2$  supports the hypersensitive transition  ${}^5D_0 \rightarrow {}^7F_2$ , indicating that europium ions are situated in a polarizable chemical environment [9, 29-31].

which is achieved in all  ${}^5D_0 \rightarrow {}^7F_2$  transitions. The Au-0.2 glass exhibits the highest values for emission cross-section, optical gain, and gain bandwidth [27]. The electric dipole transition probability ( $A_{ed}$ ) is given by equation (2)

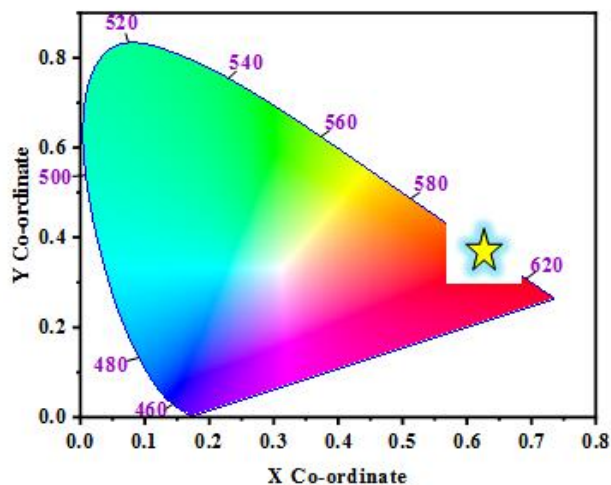
### Decay curve, quantum efficiency and CIE chromaticity analysis

The experimental decay time ( $\tau_{exp}$ ) is described by the equation  $I = I_0 \exp(-t/\tau)$ . The measured decay times ( $\tau_{exp} = 1335.1, 1773, 2044.9 \mu s$ ) are shorter than the calculated radiative lifetimes ( $\tau_{cal} = 1558, 1895.5, 2000.1 \mu s$ ). This discrepancy arises from the differences in radiative and non-radiative decay rates. The quantum efficiency ( $\eta$ ) of the corresponding levels is calculated as the ratio of the experimental lifetime to the evaluated lifetime, expressed as a percentage: 93.5%, 85.7%, and 97.9%, respectively. Figure 4 depicts the calculated CIE chromaticity coordinates for the synthesized glasses, with values falling within the NTSC threshold for

red phosphor ( $x = 0.67, y = 0.33$ ). This suggests the glasses' as ideal candidate for red laser applications.

**Table 4.** The radiative properties of the synthesized glasses

| Transition                | Parameters and units   | Au-0.0 | Au-0.2 | Au-0.5 | Au-0.8 | Au-1.0 |
|---------------------------|--|--------|--------|--------|--------|--------|
| $^5D_0 \rightarrow ^7F_1$ | $\lambda_p, \text{nm}$   | 592    | 590    | 591    | 591    | 592    |
|                           | $A_{\text{rad}}, \text{s}^{-1}$  | 180    | 180    | 182    | 183    | 183    |
|                           | $\beta_r, \%$  | 32.19  | 30.9   | 29.5   | 33.6   | 30.1   |
|                           | $\sigma, (\times 10^{-22} \text{cm}^2)$  | 3.067  | 4.047  | 4.23   | 3.78   | 3.1    |
|                           | $\sigma \times \text{FWHM}$<br>( $\times 10^{-20} \text{cm}^2$ )                     | 0.611  | 0.610  | 0.610  | 0.610  | 0.617  |
|                           | $\sigma \times \tau_{\text{cal}}$<br>( $\times 10^{-25} \text{cm}^2 \text{s}^{-1}$ ) | 5.431  | 6.87   | 6.85   | 6.94   | 5.1    |
| $^5D_0 \rightarrow ^7F_2$ | $\lambda_p, \text{nm}$   | 614    | 613    | 614    | 614    | 614    |
|                           | $A_{\text{rad}}, \text{s}^{-1}$  | 319    | 350    | 377    | 311    | 375    |
|                           | $\beta_r, \%$  | 56.4   | 59.4   | 61     | 57.2   | 61.6   |
|                           | $\sigma (\times 10^{-22} \text{cm}^2)$   | 8.113  | 7.465  | 8.48   | 6.35   | 9.55   |
|                           | $\sigma \times \text{FWHM}$<br>( $\times 10^{-20} \text{cm}^2$ )                     | 1.245  | 1.33   | 1.47   | 1.22   | 1.47   |
|                           | $\sigma \times \tau_{\text{cal}}$<br>( $\times 10^{-25} \text{cm}^2 \text{s}^{-1}$ ) | 14.367 | 12.7   | 13.7   | 11.7   | 15.7   |
| $^5D_0 \rightarrow ^7F_4$ | $\lambda_p, \text{nm}$   | 700    | 699    | 700    | 700    | 699    |
|                           | $A_{\text{rad}}, \text{s}^{-1}$  | 65.094 | 56.9   | 58.6   | 49.7   | 50.2   |
|                           | $\beta_r, \%$  | 11.5   | 9.65   | 9.49   | 9.16   | 8.27   |
|                           | $\sigma (\times 10^{-22} \text{cm}^2)$   | 3.104  | 3.06   | 3.28   | 2.76   | 2.39   |
|                           | $\sigma \times \text{FWHM}$<br>( $\times 10^{-20} \text{cm}^2$ )                     | 0.430  | 0.374  | 0.388  | 0.329  | 0.323  |
|                           | $\sigma \times \tau_{\text{cal}}$<br>( $\times 10^{-25} \text{cm}^2 \text{s}^{-1}$ ) | 5.497  | 5.19   | 5.31   | 5.07   | 3.94   |
|                           | $A_T, \text{s}^{-1}$   | 564.68 | 588.9  | 617.6  | 543.7  | 608.2  |
|                           | $\tau_{\text{cal}}, \text{ms}$   | 1.771  | 1.698  | 1.619  | 1.839  | 1.644  |



**Figure 3:** The CIE chromaticity diagram of the synthesized glasses.

### CONCLUSION

In this study, the impact of co-embedded Gold (Au) nanoparticles on alumino-borophosphate has been explored by focusing on Judd-Ofelt parameters and the radiative properties of  $\text{Eu}^{3+}$  ions. A consistent trend of  $\Omega_2 > \Omega_4 > \Omega_6$  was observed across all the glasses studied, with values ranging from  $2.05 \times 10^{-20} \text{ cm}^2$  to  $7.55 \times 10^{-20} \text{ cm}^2$ . Radiative parameters, including emission cross-section and branching ratio, were evaluated, with the  ${}^5\text{D}_0 \rightarrow {}^7\text{F}_2$  transition showing the highest branching ratio at 61.6%. This exceptional radiative emission observed suggest that these glasses are highly suitable as active media for red lasers. HRTEM analysis confirmed the presence of AuNPs with an average size of 27.65 nm. The core-shell structure of the bimetallic nanoparticles provides superior optical properties compared to monometallic nanoparticles, making them ideal for optically driven applications. The large, stimulated emission cross-section achieved highlights the potential of these materials for use in photonic devices and as active media for visible red lasers. The synthesized Au-0.30 glass sample demonstrated a high quantum efficiency ( $\eta =$

97.9%), significantly enhancing the photoluminescence intensity of the red peak, making it an excellent candidate for red laser applications.

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